#### Canard Behavior in Rate Induced Tipping

Jonathan Hahn, November 1, 2016

General framework:

$$\frac{df}{dt} = f(x, \mu, \lambda(rt))$$

# *x* is the state vector, $\mu$ is a vector of parameters, $\lambda$ is a continuous function, *r* is the rate of the forcing

For all values of  $\lambda$ , there is a stable equilibrium  $\tilde{x}(\lambda)$ 

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For all values of  $\lambda$ , there is a stable equilibrium  $\tilde{x}(\lambda)$ 

# For $r \in [0, r_c)$ , $\lambda$ changes slowly enough that if x(0) is within some neighborhood of $\tilde{x}(\lambda(0))$ , then x(t) is within some other neighborhood of $\tilde{x}(\lambda(rt))$ for all t.

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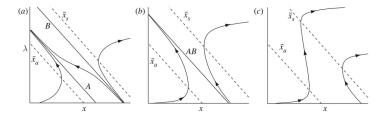
We say the state tracks the QSE for  $0 < r < r_c$ .

For  $r > r_c$ , x(t) no longer stays within the required neighborhood of the QSE. The system "tips" and we say it has *rate-dependent tipping*.

The neighborhood of the QSE that describes the tipping point can be chosen in many ways: by a given distance *R*, by the state leaving a basin of attraction of the QSE, by something else topological in the system, or by some other arbitrary choice.

# 

$$\frac{dx}{dt} = (x + \lambda)^2 - \mu$$
$$\frac{d\lambda}{dt} = r$$



# Rate-induced tipping example

$$\frac{dx}{dt} = (x+\lambda)^2 - \mu$$
$$\frac{d\lambda}{dt} = r$$

Co-moving system: set  $w = x + \lambda$ .

$$\frac{dw}{dt} = w^2 - \mu + r$$

Equilibrium at  $w = \pm \sqrt{\mu - r}$  if  $r < \mu$ .

Tipping condition:  $r > r_c$  where

$$r_c = \begin{cases} \mu - (\lambda_0 + x_0)^2 & \text{if } - x_0 < \lambda_0 < -x_0 + \sqrt{\mu} \\ \mu & \text{if } \lambda_0 \le -x_0 \end{cases}$$

# FAST-SLOW SYSTEM

QSE near a locally folded critical manifold.

$$\epsilon \frac{dx}{dt} = y + \lambda + x(x - 1)$$
$$\frac{dy}{dt} = -\sum_{n=1}^{N} x^{n}$$

 $N \ge 5$ , odd. (0,  $-\lambda$ ) is globally asymptotically stable.

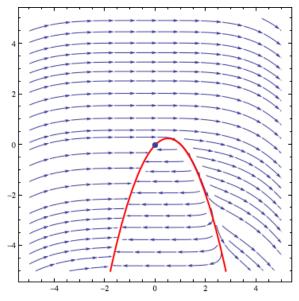
# FAST-SLOW SYSTEM

$$\epsilon \frac{dx}{dt} = y + \lambda + x(x - 1)$$
$$\frac{dy}{dt} = -\sum_{n=1}^{N} x^{n}$$

Set  $\epsilon = 0$  to find the slow manifold:  $0 = y + \lambda + x(x - 1)$ 

$$S(\lambda) = \{(x, y) \in \mathbb{R}^2 : y = -\lambda - x(x - 1)\}$$

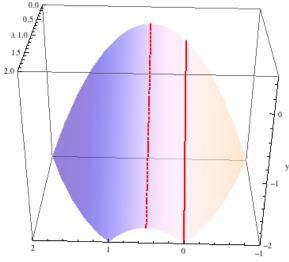
# FAST-SLOW SYSTEM



# FAST SLOW SYSTEM: RATE TIPPING

$$\epsilon \frac{dx}{dt} = y + \lambda + x(x - 1)$$
$$\frac{dy}{dt} = -\sum_{n=1}^{N} x^{n}$$
$$\frac{d\lambda}{dt} = r$$

# Critical Manifold



## Projected Reduced System

Set  $\epsilon = 0$  and differentiate the resulting equation with respect to *t* to find a system approximating the slow dynamics.

$$0 = \frac{dy}{dt} + \frac{d\lambda}{dt} + (2x - 1)\frac{dx}{dt}$$
$$\frac{dx}{dt} = \left(\sum_{n=1}^{N} x^n - r\right)(2x - 1)^{-1}$$
$$\frac{d\lambda}{dt} = r$$

#### Desingularized system

Rescale time:  $\frac{dt}{d\tau} = -(2x - 1)$ 

$$\frac{dx}{d\tau} = \left(r - \sum_{n=1}^{N} x^n\right)$$
$$\frac{d\lambda}{d\tau} = -r(2x - 1)$$

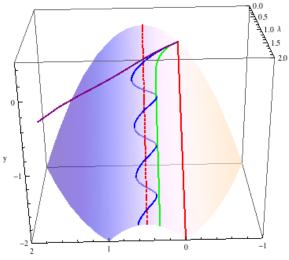
This reverses the direction of time on the repelling part of the critical manifold.

For  $0 < r < \sum_{n=1}^{N} (1/2)^n$ , all trajectories within the attracting part of the critical manifold converge to  $x^*$  where  $r = \sum_{n=1}^{N} x^*$ .

# Critical rate

$$r_c = \sum_{n=1}^N (1/2)^n$$

# Solutions



# CO-MOVING SYSTEM

$$\epsilon \frac{dx}{dt} = y + \lambda + x(x - 1)$$
$$\frac{dy}{dt} = -\sum_{n=1}^{N} x^{n}$$
$$\frac{d\lambda}{dt} = r$$

Create a co-moving system:  $w = y + \lambda$ 

$$\epsilon \frac{dx}{dt} = w + x(x-1)$$
$$\frac{dw}{dt} = r - \sum_{n=1}^{N} x^{n}$$

# CO-MOVING SYSTEM

$$\epsilon \frac{dx}{dt} = w + x(x - 1)$$
$$\frac{dw}{dt} = r - \sum_{n=1}^{N} x^{n}$$

The equilibrium  $(x^*, w^*)$  in the co-moving system is given by the solution to

$$\sum_{n=1}^{N} (x^*)^n = r$$
$$w^* = -(x^*)^2 + x^*$$

# Hopf bifurcation analysis

The Jacobian at this equilibrium is:

$$\begin{bmatrix} (2x_1^* - 1)/\epsilon & 1/\epsilon \\ \sum_{n=1}^N -n(x_1^*)^{n-1} & 0 \end{bmatrix}$$

and the eigenvalues of the Jacobian are

$$\frac{2x_1^* - 1 \pm \sqrt{(1 - 2x_1^*)^2 - 4\epsilon \sum_{n=1}^N n(x_1^*)^{n-1}}}{2\epsilon}$$

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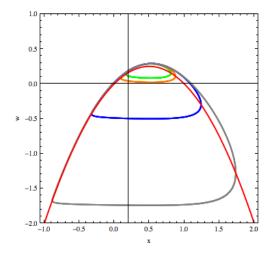
### Hopf bifurcation analysis

#### Eigenvalues:

$$\frac{2x_1^* - 1 \pm \sqrt{(1 - 2x_1^*)^2 - 4\epsilon \sum_{n=1}^N n(x_1^*)^{n-1}}}{2\epsilon}.$$

When  $x^* < 1/2$ , the equilibrium is stable, which is in agreement with the previous conclusion that the system does not tip for  $r < \sum_{n=1}^{N} (1/2)^n = r_c$ . As *r* increases, so does  $x^*$ , so when  $r = r_c$ , and  $x^* = 1/2$  the pair of eigenvalues cross the imaginary axis, and a Hopf bifurcation occurs.

Periodic orbits in co-moving system



Periodic orbit expands rapidly as in a canard explosion.

# VAN DER POL OSCILLATOR

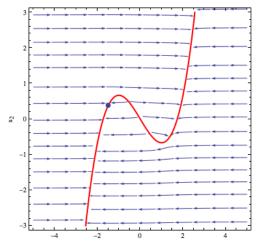
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$$\frac{dx_2}{dt} = -x_1 - \alpha$$

# VAN DER POL OSCILLATOR

$$\epsilon \frac{dx}{dt} = x_2 + (x_1 - \frac{x_1^3}{3}) + \lambda$$
$$\frac{dx_2}{dt} = -x_1 - \alpha$$
$$\frac{d\lambda}{dt} = r > 0$$

 $\alpha>1$  is fixed, and a stable equilibrium exists at  $(-\alpha,-\alpha+\frac{\alpha^3}{3}-\lambda)$ 

Phase portrait for  $\lambda = 0$ .



**x**1

# Co-moving van der Pol System

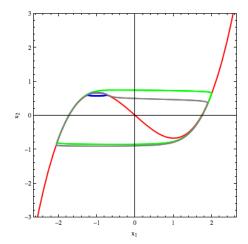
Set  $w = x_2 + \lambda$ :

$$\epsilon \frac{dx_1}{dt} = w + (x_1 - \frac{x_1^3}{3})$$
$$\frac{dw}{dt} = -x_1 - \alpha + r$$

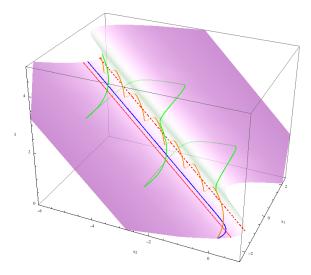
This is the classic van der Pol system!

# CANARD PROGRESSION

As *r* increases beyond  $\alpha - 1$ , there is a canard explosion.



# Orbits in original system





Does spiraling behavior still count as "tracking"?

If so, is the critical rate for spiraling really a "tipping point"?

How do we prove this sprialing occurs when we can't reduce to a "co-moving system"?